

An Efficient ADER-DG Scheme for Simulation of Seismic Waves in Poroelastic Media

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Summary

The equation of motion for poroelastic media based on Biot's theory contains a stiff reactive term that poses a key computational challenge for numerical methods. We present an efficient ADER-DG scheme (discontinuous Galerkin scheme with Arbitrary high-order DERivative time stepping) for solving problems with the stiff term utilizing a novel block-wise back-substitution algorithm. For polynomials of degree 6, the number of floating-point operations is reduced by a factor of 25 compared to the standard LU decomposition. Additionally, the block-wise back-substitution is mapped to a sequence of small matrix-matrix multiplications, for which code generators are available to generate highly optimized code. This way, we achieve excellent node-level performance. We verify the new scheme thoroughly against (semi)analytical or numerical reference solutions in canonical problems of increasing complexity using a point source. Additionally, by utilizing a clustered local time-stepping scheme, time to solution is reduced by a factor of 6 to 10 compared to the global time stepping. Performance analysis on the SuperMUC-NG supercomputer demonstrates high computational efficiency of our implementation and its potential for extreme-scale simulations.

ADER-DG Scheme

Governing equation for wave propagation in a poroelastic medium in matrix-vector form:

$$\frac{\partial q}{\partial t} + A \frac{\partial q}{\partial x} + B \frac{\partial q}{\partial y} + C \frac{\partial q}{\partial z} = E \cdot q$$

q is the vector of unknowns consisting of six stress components, three particle velocity components, pore pressure and three relative fluid velocities:

$$q = (\sigma_{xx}, \sigma_{yy}, \sigma_{zz}, \sigma_{xy}, \sigma_{yz}, \sigma_{xz}, u, v, w, p, u_f, v_f, w_f)^T$$

$E \cdot q$ is the viscous dissipation term and presents the challenging stiff term

Matrices A, B, C and E contain material properties and have specific sparsity patterns (see Fig 1)

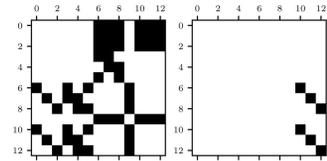


Figure 1
Sparsity pattern of matrix
Left $A + B + C$
Right E

Time integration is performed using a predictor-corrector scheme, in which the locally predicted solution is corrected using numerical fluxes. The Cauchy-Kovalevski procedure, used as the predictor for (visco-)elastic media, is unstable in the presence of the stiff term. Therefore, we adopt the space-time predictor proposed by de la Puente *et al.*, 2008. The idea is to express the solution on the reference element in space and time using a polynomial basis functions

$$Q_p(\xi, \eta, \zeta, \tau) = Q_{pls} \cdot \chi_l(\tau) \cdot \psi_s(\xi, \eta, \zeta)$$

Following the derivation by de la Puente *et al.*, 2008, the system of equations for Q_{pls} can be written in the form

$$Y_{pkrqls} Q_{pls} = R_{pkr}$$

The operator Y_{pkrqls} and the right-hand side R_{pkr} are

$$Y_{pkrqls} = \delta_{pq} W_{rs} M_{kl} - \delta_{pq} K_{rs}^i M_{kl} - E_{pq}^* S_{rs} M_{kl} + \sum_{j=1}^3 A_{pq}^{j*} S_{rs} K_{kl}^j$$

$$R_{pkr} = \delta_{pq} w_r M_{kl} Q_{ql}^0$$

where

$$W_{rs} = \chi_r(1) \chi_s(1)$$

$$w_r = \chi_r(0)$$

$$S_{rs} = \int \chi_r \chi_s d\tau$$

$$K_{rs}^\tau = \int \frac{\partial \chi_r}{\partial \tau} \chi_s d\tau$$

$$M_{kl} = \int \phi_k \phi_l dV$$

$$K_{kl}^\alpha = \int \phi_k \frac{\partial \phi_l}{\partial \alpha} dV \quad \alpha \in \{\xi, \eta, \zeta\} \text{ or } \{1, 2, 3\}$$

and matrices A^j are Jacobian matrices composed from matrices A, B and C . See Wolf *et al.*, 2022 for detailed derivation and discussion

We present an efficient ADER-DG scheme for simulating seismic waves in poroelastic media

Our novel algorithm for treating the stiff source term reduces the number of operations by a factor of 25 compared to the standard LU decomposition algorithm for basis functions of degree 6

Our implementation in SeisSol (www.seissol.org) reduces the time to solution by a factor of about 10 compared to an ideally performing LU decomposition



Wolf, Galis, Uphoff, Gabriel, Moczo, Gregor, Bader, 2022.
An Efficient ADER-DG Local Time Stepping Scheme for 3D HPC Simulation of Seismic Waves in Poroelastic Media, Journal of Computational Physics, 455, doi: 10.1016/j.jcp.2021.110886

The Novel Block-Wise Back-Substitution Algorithm

By unrolling the multi-indices pkr and qls to i and j , respectively, the system $Y_{pkrqls} Q_{pls} = R_{pkr}$ can be written in standard matrix form, $Ax = b$

Choosing k/l as the slowest and r/s as the fastest-running indices leads to a specific block-structure of the system matrix (see Fig 2).

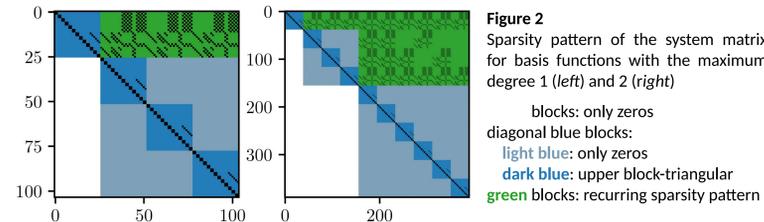


Figure 2
Sparsity pattern of the system matrix for basis functions with the maximum degree 1 (left) and 2 (right)
blocks: only zeros
diagonal blue blocks:
light blue: only zeros
dark blue: upper block-triangular
green blocks: recurring sparsity pattern

Our novel algorithm (see Fig 3) utilizes

- the block-wise back-substitution exploiting the block-triangular structure of the system matrix
- the upper triangular sparsity pattern of the matrix E^*
- the internal structure of blue and green blocks
- tensor operations

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1  $b \leftarrow Q^0 \circ (S^{-1}w)$ ;
2 for  $n \leftarrow N$  downto 0 do
3   // dark and light-blue blocks together
4    $m \leftarrow (B_{n-1}, B_n]$ ;
5   for  $p \leftarrow Q$  downto 1 do
6      $Q_{pm} \leftarrow b_{pm} \cdot (Z - E_{pp}^* D)^{-T}$ ;
7     for  $o \leftarrow 1$  to  $p-1$  do
8        $b_{om} \leftarrow b_{om} + G_{op} \cdot Q_{pm}$ ;
9     end
10  end
11 // green blocks
12 if  $n > 0$  then
13    $b \leftarrow b - A^* \times_1 Q_m \times_2 \tilde{K}_{m1}^\xi$ ;
14    $b \leftarrow b - B^* \times_1 Q_m \times_2 \tilde{K}_{m1}^\eta$ ;
15    $b \leftarrow b - C^* \times_1 Q_m \times_2 \tilde{K}_{m1}^\zeta$ ;
16 end
17 end
    
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Figure 3
The scheme of our novel algorithm (Algorithm 3 in Wolf *et al.*, 2022)

Performance

Advantages

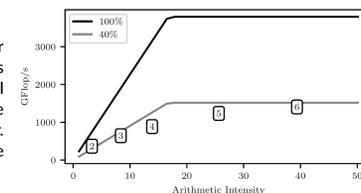
- lower number of FLOPs compared to standard LU decomposition (see Tab 1)
- lower memory requirements than standard LU decomposition (see Tab 1)
- because the algorithm is defined via tensor operations, the code generator YATeTo (Uphoff *et al.*, 2020) can be used to translate/map the tensor operations to small matrix-matrix multiplications, for which highly optimized architecture-specific code generators are available
- time to solution with our SeisSol implementation is about 10 times shorter than with ideally performing LU decomposition
- the algorithm can be combined with local time stepping - for example, for LOHp test, the time to solution is 6 to 10 times shorter than with global time stepping depending on the number of used nodes and convergence order

Table 1 Comparison of number of floating-point operations and memory requirements of standard LU decomposition (LU) and our optimized algorithm 3 for the space-time predictor approach (STP). [N - maximum degree of basis functions, n - number of unknowns]

N	n	FLOps			storage		
		LU	STP	reduction	LU [MB]	STP [MB]	reduction
4	2275	1.04×10^7	7.13×10^5	14.5	39.5	6.4×10^3	6179
5	4368	3.82×10^7	1.94×10^6	19.7	145.6	7.5×10^3	19453
6	7644	1.17×10^8	4.72×10^6	24.8	445.9	8.8×10^3	50816

Figure 4

The roofline model shows that our implementation in SeisSol performs at approx. 40% of the theoretical peak performance of a single node of SuperMUC-NG supercomputer. The numbers in squares depict the maximum degree of basis functions.



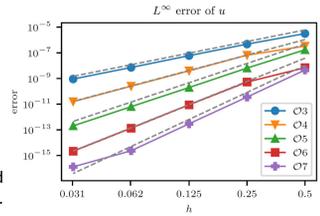
Numerical Tests - Verification

Convergence test

- planar wave as the initial condition
- analytical reference solution
- expected convergence rate achieved for all unknowns

Figure 5

Convergence plot for particle velocity u . Dashed gray lines depict the expected convergence rates.



Point source in homogenous space

- explosive point source at (0,0,0)
- semi-analytical reference solution (Karpfinger *et al.*, 2009)
- excellent agreement for inviscid as well as viscous cases

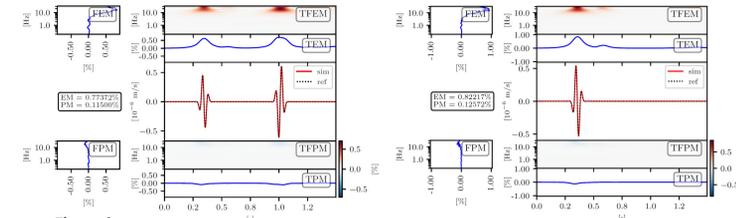


Figure 6

Detailed misfit plots for particle velocity u at location (575m, 575m, 575m) for materials with the pore space filled by inviscid fluid (left) and viscous fluid (right).

Boundary conditions

- explosive point source
- semi-analytical reference solution (gar6more3D by Diaz *et al.*, 2008)
- only inviscid case (due to limitations of the reference solution)
- excellent agreement for free surface boundary condition and internal interface

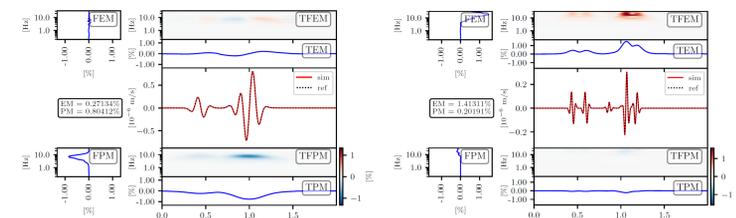


Figure 7

Detailed misfit plots for particle velocity u in inviscid case. Left: Free-surface boundary condition - receiver at a shallow depth (due to limitations of the reference solution). Right: Internal interface - receiver 500m behind the interface captures transmitted waves.

Layer over halfspace (LOHp)

- explosive line source
- 2D finite difference reference solution (Gregor *et al.*, 2021)
- important for most practical applications:
 - excellent agreement for solid particle velocities with standard resolution
- important for specific applications:
 - if the relative fluid velocities must be calculated at or very close to the free surface or an internal interface, finer resolution is necessary in both methods

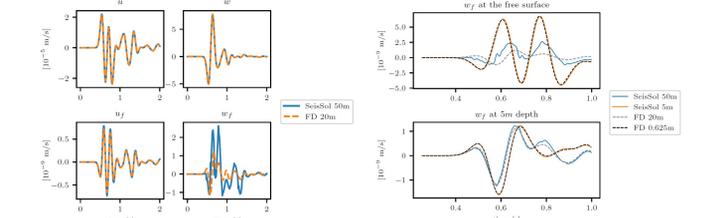


Figure 8

Comparison of solid particle velocities (u and w) and relative fluid velocities (u_f and w_f) in the LOHp model. Left: The receiver at the free surface shows excellent agreement for solid particle velocities using standard resolution. Right: Illustration of the small-scale effect of the diffusive slow P-wave near the free surface using fine resolution.

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