A Three-Dimensional Nonlinear Finite Element Formulation for Geometrically Exact Kirchhoff Rods

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Introduction

Geometrically exact beam models

- Beams as structural models appropriate to describe the mechanics of bodies, whose length dimension is much larger than the transverse dimensions
- Geometrically exact (GE), if deformation and stress measures are work-paired
- Deduction of beam models from the 3D continuum:



Reference Triad Field

- Choice of reference triad as crucial step, that influences factors like continuity • requirements, objectivity, convergence behaviour, robustness and complexity
- SR mapping defines triad $g_{SR,i}$ out of given triad \bar{g}_i and given vector $g_{SR,i}$:

$$\mathbf{g}_{SR,j} = \bar{\mathbf{g}}_j - \frac{\bar{\mathbf{g}}_j \cdot \mathbf{g}_{SR,1}}{1 + \bar{\mathbf{g}}_1 \cdot \mathbf{g}_{SR,1}} \left(\mathbf{g}_{SR,1} + \bar{\mathbf{g}}_1 \right) \quad \text{for} \quad j = 2,3$$

- Direct application of the SR mapping leads to non-objectivity (SR element)
- An objective formulation follows with the procedure (NSRISR element): 1) Define reference triads at nodal positions ξ_1^n and time t_k via SR mapping: $\mathbf{g}_{SR,j} = \mathbf{g}_{ref,j}(\xi_1^n, t_k), \ \bar{\mathbf{g}}_i = \mathbf{g}_{ref,i}(\xi_1^n, t_{k-1}), \ \mathbf{g}_{SR,1} = \mathbf{g}_1(\xi_1^n, t_k)$
- Constrained position vector due to Bernoulli assumption of rigid cross sections: $\mathbf{X}(\xi_1,\xi_2,\xi_3) = \mathbf{R}(\xi_1) + \xi_2 \mathbf{G}_2(\xi_1) + \xi_3 \mathbf{G}_3(\xi_1)$ if $h \ll l$ $\mathbf{x}(\xi_1,\xi_2,\xi_3) = \mathbf{r}(\xi_1) + \xi_2 \mathbf{g}_2(\xi_1) + \xi_3 \mathbf{g}_3(\xi_1)$
- Model reduction from 3D Boltzmann continuum to 1D Cosserat continuum: 3 DoF: $\mathbf{x}(\boldsymbol{\xi}) \rightarrow 6$ DoF: $\mathbf{r}(\xi_1)$, $\mathbf{q}(\xi_1)$ with $\mathbf{g}_i = \Lambda(\mathbf{q}(\xi_1))\mathbf{e}_i$ for i = 1, 2, 3
- Reduction from Reissner theory to Kirchhoff theory possible for slender beams: 6 DoF: $r(\xi_1), q(\xi_1) \rightarrow 4$ DoF: $r(\xi_1), \varphi(\xi_1)$ since $g_1 \times r' \doteq 0$

Aim

- Existing GE FEM beam formulations exclusively of Reissner type
- Development of a GE FEM beam formulation according to Kirchhoff theory

Governing Equations

Kinematics

2) Define interior reference triad field at position ξ_1 and time t_k via SR mapping: $\mathbf{g}_{SR,j} = \mathbf{g}_{ref,j}(\xi_1, t_k), \ \bar{\mathbf{g}}_i = \mathbf{g}_{ref,i}(\xi_1^n, t_k), \ \mathbf{g}_{SR,1} = \mathbf{g}_1(\xi_1, t_k)$

Numerical Examples

Inversion of an elastic helix

- Clamped helix: E=2G=1, $I=10^3$, with h=10, R=H=I/($9\pi\sqrt{2}$)
- The external moment $M=18\pi EI_2/I$ exactly inverts the elastic helix





Initial (red) and deformed (green) geometries

Instability of an elastic ring



Description of the beam centerline and the material frame ${f g}_i$ via the 4 DoF ${f r}$, arphi

Introduction of a reference triad field, which is part of the solution:

 $g_{ref,1} = r'/||r'||, \quad g_{ref,2} = f_2(r), \quad g_{ref,3} = f_3(r)$

- Determination of the material frame via a relative rotation of reference frame: $\mathbf{g}_1 = \mathbf{g}_{ref,1}, \, \mathbf{g}_2 = \mathbf{g}_{ref,2} \cos \varphi + \mathbf{g}_{ref,3} \sin \varphi, \, \mathbf{g}_3 = \mathbf{g}_{ref,3} \cos \varphi - \mathbf{g}_{ref,2} \sin \varphi$
- Derivation of geometrically exact spatial deformation measures:

$$\omega = \begin{pmatrix} \tau + \varphi' - \tau_0 - \varphi'_0 \\ (g_2 \cdot \kappa) - (G_2 \cdot \kappa_0) \\ (g_3 \cdot \kappa) - (G_3 \cdot \kappa_0) \end{pmatrix}_{g_i} \text{ and } \epsilon = ||\mathbf{r}'|| - \epsilon$$

The torsion of the reference system and the centerline curvature are:

$$au = \mathbf{f}_2'(\mathbf{r}) \cdot \mathbf{f}_3(\mathbf{r})$$
 and $\kappa = rac{\mathbf{r}' imes \mathbf{r}''}{||\mathbf{r}'||^2}$

Constitutive law

Derivation of spatial stress resultants from hyperelastic stored energy function:

$$\mathbf{n} = n\mathbf{g}_1 \quad \text{with} \quad n = \frac{\partial_{rel} \Pi_{int}(\epsilon, \omega)}{\partial \epsilon} \quad \text{and} \quad \mathbf{m} = \frac{\partial_{rel} \Pi_{int}(\epsilon, \omega)}{\partial \omega}$$

e.g. $\Pi_{int}(\epsilon, \omega) = \frac{1}{2} EA\epsilon^2 + \frac{1}{2} \omega^T \mathbf{c}_m \omega \quad \text{with} \quad \mathbf{c}_m = \text{diag}(GI_P, EI_2, EI_3)_{\mathbf{g}_i}$

Weak form

$$\int_{0}^{L} \left[\left(\delta \alpha \mathbf{t} + \frac{\mathbf{r}'}{|\mathbf{r}'|^{2}} \times \delta \mathbf{r}' \right)' \cdot \mathbf{m} + \delta \mathbf{r}' \cdot \mathbf{n} - \delta \mathbf{r} \cdot \tilde{\mathbf{f}} - \left(\delta \alpha \mathbf{t} + \frac{\mathbf{r}'}{|\mathbf{r}'|^{2}} \times \delta \mathbf{r}' \right) \cdot \tilde{\mathbf{m}} \right] ds$$

Conclusions

Geometrically exact Kirchhoff beam element

- Strong enforcement of Kirchhoff constraint
- Kirchhoff kinematics based on objective reference triad interpolation

$$-\left[\delta \mathbf{r} \cdot \mathbf{f}\right]_{\Gamma_{\sigma}} - \left[\left(\delta \alpha \mathbf{t} + \frac{\mathbf{r}'}{|\mathbf{r}'|^2} \times \delta \mathbf{r}'\right) \cdot \mathbf{m}\right]_{\Gamma_{\sigma}} = 0$$

with external line loads $\tilde{\mathbf{f}}, \tilde{\mathbf{m}}$ and primary variables $\mathbf{r}, \delta \mathbf{r}, \varphi, \delta \alpha$

Discretization

Discretization of beam centerline with third order Hermite shape functions:

$$\mathbf{r}_{h}(\xi) = \sum_{i=1}^{2} H_{d}^{i}(\xi) \mathbf{d}^{i} + \frac{l_{ele}}{2} \sum_{i=1}^{2} H_{t}^{i}(\xi) \mathbf{t}^{i}$$

with nodal positions d^i , nodal tangents t^i and the element length l_{ele} \rightarrow Fulfilment of completeness, interpolation property, C¹-continuity at nodes Discretization of relative angle with third order Lagrange shape functions:

$\varphi_h(\xi) = \sum L^i(\xi) \hat{\varphi}_i$

- C¹-continuity of beam geometry through Hermite interpolation
- Efficient model reduction possible for isotropic and torsion free applications
- Objectivity, path independence, consistent convergence order and locking avoidance theoretically predicted and numerically confirmed

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