



## **Conceptual Framework of Energy Dissipation During Disintegration in Rock Avalanches**

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Rock avalanches usually progress through three consecutive phases: Detachment (Phase 1), Disintegration (Phase 2), and Flow (Phase 3). While significant advances have been achieved in modeling Rock Avalanche Phase 1 (Detachment) and Phase 3 (Flow), the crucial link between both during Phase 2 (Disintegration) is still poorly understood. Disintegration of the detached rock mass is often initiated as soon as sliding starts, and in situ measurements are impossible due to the excessive energy release equivalent to multiple nuclear explosions. Better understanding the energy dissipation during Phase 2, and the resulting residual kinetic energy that propels the rock avalanche in Phase 3, is one of the keys to defining the mechanical properties of the avalanche in the runout zone and thus also the resisting force within the avalanche. This paper is a review of our knowledge of energy dissipation in rock avalanches with a focus on processes like friction, collision, fragmentation, comminution, entrainment and explosion during the phase of disintegration. We distinguish between energy sources and sinks and consider not only physical processes, but also chemical alterations that might occur at high temperatures. With that, we make a contribution to improve our understanding of Phase 2 "Disintegration," which is needed for accurately modeling rock avalanches and assessing their hazard potential.

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## INTRODUCTION

Rock avalanches are defined as "extremely rapid, massive, flow-like motion of fragmented rock from a large rock slide or rock fall" (Hungr et al., 2014). Due to their high velocity, volume and runout distance, rock avalanches have a significant impact on human activities in mountain areas, can seriously damage infrastructure and settlements and can cause high numbers of casualties (Evans et al., 2006; Legros, 2006; Hewitt et al., 2008). Landslides resulting from large-scale rock-slope failures are especially hazardous; in the 20th century, disasters of this type have killed more than 50,000 people globally (Evans et al., 2006). As a consequence of increasing population density and the development of infrastructure in mountain areas, the number of elements at risk is growing and accelerating the vulnerability to landslide hazards (Fischer, 1999; Korup, 2005; Hungr, 2006; Legros, 2006). At the same time, the number of massive rock failures from permafrost warming appears to be increasing with potentially disastrous consequences especially when causing rock-ice avalanches with high mobility (Haeberli et al., 2004; Huggel, 2009; Huggel et al., 2012; Krautblatter et al., 2013; Krautblatter and Leith, 2015) or causing flooding after impacting lakes (Haeberli et al., 2016; Knapp et al., 2018).

Better understanding the disintegration (Phase 2; Figure 1) is key to defining mechanical properties like grain size composition and content of large blocks in the runout zone and therefore the hazard potential of rock avalanches. Current approaches based on Mohr-Coulomb friction models adequately describe the detachment processes (Phase 1; Figure 1) and its energy dissipation (Maddock, 1986) or the rock avalanche flow (Phase 3; Figure 1) utilizing fluid (Bingham) or snow avalanche (Voellmy) analogs with adequate parameterization (Hungr, 2006; Christen et al., 2010; Preuth et al., 2010; Pudasaini and Krautblatter, 2014; Pudasaini and Mergili, 2019). For Phase 2, some models on dynamic fragmentation were just developed (e.g., Zhao et al., 2017; Ghaffari et al., 2019), whereas other disintegration processes, e.g., heat transfer and phase transitions still represent major research gaps. This situation is mainly related to insufficient understanding of energy dissipation during Phase 2, and the resulting residual kinetic energy that propels the ensuing rock avalanche (Phase 3). The material properties of the avalanche result from these energetic processes and from the material being overrun. Only by understanding disintegration, will more precise modeling of rock avalanches and their hazard potential be possible. In this paper, we are going to primarily concentrate on the intrinsic properties of rock avalanches that influence disintegration, and we are going to focus on disintegration and energy dissipation in Phase 2, that is directly after the detachment.

# ENERGY DISSIPATION DURING DISINTEGRATION

#### **Disintegration Processes**

Large rock-slope failures usually undergo different stages of downhill movement which may occur consecutively (Abele, 1974): (i) The rock mass moves as a coherent block, and translational shearing occurs along the contact of the bottom of the rock avalanche and the ground surface. (ii) Subsequently, differential movement of individual blocks initiates crushing of the original rock mass. (iii) If the coherent rock mass loses its internal cohesion and disintegrates intensely (shattering) it can evolve into a rapid granular flow (Pollet and Schneider, 2004), which is defined as the distributed shear motion of a group of clasts where individual grains interact with each other and with the boundaries of the moving flow (Dufresne and Davies, 2009). The result can be a highly fragmented (pulverized) rock mass which consists of angular grains of all sizes down to <1  $\mu$ m (**Figures 2A–D**; e.g., Davies and McSaveney, 2012).

To decipher individual processes during disintegration, two types of disintegration can be distinguished: (i) *static disintegration*, a collision-free process driven mainly by gravity, and (ii) *dynamic disintegration*, referring to particle comminution by grain-to-grain collisions driven by motion. Disintegration refers to fracturing by rapid changes in stress coupled with sudden (un-)loading caused by bending, transverse shearing or delamination of the rock mass creating large blocks, thin vertical slabs or, horizontal sheets, respectively (Erismann and Abele, 2001). Static disintegration is an essential precursor for dynamic disintegration as it creates fractures along which further relative shearing and fragmentation can occur. Shearing along predefined bedding and foliation planes induces shear crushing and the creation of a granular layer.

## **Energy Sources and Sinks**

Energy dissipation in rock avalanches occurs by transformation of the total energy into thermal energy, acoustic energy or inelastic deformation energy (Nicoletti and Sorriso-Valvo, 1991), where due to the law of energy conservation, the final energy available for mechanical work is less than the initial amount. The energy release is often in the range of dozens to more than a thousand Hiroshima bombs (~15 kt TNT or 63 TJ each) for large rock avalanches. Recent work also emphasizes the energy transfer into chemical reactions and phase transitions (Anders et al., 2010; Mitchell et al., 2015). Energy "release" and "consumption" describe the transfer of energy into a different form. Energy in rock avalanches is released by friction, collision and fracturing. Far from a continuous process, energy release is concentrated at points of impact with the ground surface and obstacles where major friction and disintegration of the rock mass is initiated (Erismann and Abele, 2001).

Field conditions constraining energy dissipation can be derived from (i) paleotopography (Nicoletti and Sorriso-Valvo, 1991), (ii) compressive and extensional flow structures in the rock-avalanche deposits (Hewitt, 2006; Dufresne and Davies, 2009; Dufresne et al., 2015), (iii) positions inside the flow recording differences between intact rock and major shear zones (Pollet et al., 2005), (iv) the sedimentological record (Yarnold, 1993; Weidinger et al., 2014) with (v) fine-sediment signatures (Reznichenko et al., 2012), and (vi) melting mineral formation (Weidinger and Korup, 2009). Referring to (i), Nicoletti and Sorriso-Valvo (1991) differentiate dissipation types and rates dependent on geomorphic controls along the runout path: The low-energy dissipative type refers to rock avalanches which are, for example, channelized in narrow valleys. Here, little potential energy is dissipated to other processes than kinetic energy, and mobility is enhanced. The *moderate-energy dissipative type* refers to radial spreading "free from lateral constraints," resulting in moderate mobility. Finally, the high-energy dissipative type describes running across a narrow valley and impacting against the opposite, at best perpendicular slope, which results in low mobility. Here, most initial energy is dissipated to energy sinks, and only little is left for the transfer into kinetic energy.

During disintegration,  $\sim 20-50\%$  of the potential energy is consumed (Locat et al., 2006; Haug et al., 2016). Considering multiple energy sinks in **Figure 1**, A–D (friction, inelastic collision, entrainment, and crustal deformation) cause heating to some degree, E and F (chemical energy consumption and phase transition) require latent energy for phase transitions, and G– I (dust production and bouncing, sound and microseismicity, and momentum exchange) act to export energy outside the impact/disintegration zone. The relative importance of D, G, H, and I (compression, dust production and bouncing, sound and microseismicity, and momentum exchange) has yet to be determined, but Erismann and Abele (2001) assumed that they are of minor importance. If A–C (friction, inelastic collision,



and entrainment) have a major share in the energy dissipation and cause a mean frictional shear resistance whereas E-G(chemical energy consumption, phase transformation, and dust production) consume energy for phase transition, the rate of frictional heat generation per unit area *Q* is

$$\dot{Q} = \tau v - \phi = \mu_k \sigma_n v - \phi \tag{1}$$

where  $\tau$  is the average frictional shear resistance, v is the average velocity,  $\phi$  is the heat sink-rate due to latent heat,  $\mu_k$  is the kinetic coefficient of friction and  $\sigma_n$  is the normal stress across the sliding plane (Maddock, 1986). Effective latent heat sinks could be from decarbonation of dolomite and calcite in sedimentary rock failures (Mitchell et al., 2015) or from phase transitions of water during melting and vaporization (De Blasio and Medici, 2017). The heat flow away from a source (e.g., a sliding plane) can be calculated by 1D-heat diffusion (Carslaw and Jaeger, 1959; Mitchell et al., 2015), where the temperature increase  $\Delta T$  within the observed slip zone is

$$\Delta T(x,t) = \frac{1}{2\rho c \sqrt{\kappa \pi}} \int_{0}^{t} \frac{\tau(t')\nu(t') - \phi(t')}{\sqrt{t - t'}} e^{\frac{-x^2}{4\kappa(t - t')}} dt'$$
(2)

where x is the distance from the slip zone, t is time,  $\rho$  is mass density,  $\kappa$  is thermal diffusivity, c is heat capacity.

## Physical Processes

#### Fragmentation/Collision/Comminution

Fragmentation describes the initiation and propagation of fractures and breaking apart and movement of grains (Turcotte, 1997). The related process energy is both linked to the length of the crack extension within existing grains (microcracking) and to the surface energy of the new created grains during comminution (Bieniawski, 1967; Hamdi et al., 2008). Fragmentation occurs as a *static* (Eberhardt et al., 2004; Wang et al., 2011; Zhang, 2016) or *dynamic* process (Pollet and Schneider, 2004; Crosta et al., 2007; Imre et al., 2010; Zhang et al., 2019). *Static fracture* occurs before any collision triggers the disintegration of a mass, whereas *dynamic fragmentation* shows more intense disintegration, e.g., in shear zones at the base of rock avalanches.

Grains fragment quickly under high local pressures and, thus, general intergranular effective stress and the frictional resistance to shear are reduced (Bowman et al., 2012). In laboratory experiments, the overburden strain-rate is directly related to the fragmentation process. If load is applied sufficiently quickly, particles will dynamically fragment and the kinetic energy of the resulting fragments will cause collisions with surrounding particles. Under dynamic disintegration, kinetic energy is dispersed through the system as colliding particles undergo further fragmentation (Rait and Bowman, 2010). The higher the spatial concentration of simultaneously-fragmenting grains, the lower the effective direct stress on the grain flow (Davies and McSaveney, 2009). Thereby, the basal sliding friction dissipates upward and laterally through the mass, which causes the slabs at the bottom to come to rest first. Thus, slabs higher in the moving mass travel further than the ones lower down (Erismann and Abele, 2001; Pollet and Schneider, 2004). Grainto-grain collisions require an unconfined environment, in which particles can move freely. In such a case the highest levels of friction, crushing and collision occur in the lower part of a rock





avalanche due to high compressing forces and large differences in velocity between the moving particles and the ground (Erismann and Abele, 2001). Running on dry rock substrates, it is mainly fragmentation that leads to an increased travel length of the rock avalanche (Pollet and Schneider, 2004; McSaveney and Davies, 2007; Davies and McSaveney, 2009). After Haug et al. (2016), increased fragmentation mostly affects the front of a rock avalanche traveling further, whereas the center of mass crucial for energy considerations, is hardly displaced or decelerates. In comparison to previous papers stating that fragmentation accelerates the flow (Bowman et al., 2012; Langlois et al., 2015), Haug et al. (2016) confirmed that high fragmentation rather favors a more energy-efficient transport mode yielding longer runouts without acceleration.

In either case, fragmentation is considered an "energy sink" (Locat et al., 2006; Crosta et al., 2007; Haug et al., 2016). Haug et al. (2016) propose that static fragmentation may use up to 50% of the potential energy. Also, Ghaffari et al. (2019) postulate that the kinetic energy is only a small portion of dissipated energy during fragmentation, and the energy rather

transfers into intergranular collision and friction. Thereby, it is important to note that the energy input for grain-internal "microcracking" weighs far more than for "macro-fragmentation," i.e., the formation of new grains (Ouchterlony et al., 2004; Hamdi et al., 2008). Zhao et al. (2017) quantify the energy dissipation by friction and plastic deformation to ~90%, and the energy needed by bond breakage to <5%. Plus, the smaller the grain size becomes, the more energy is needed for comminution (Locat et al., 2006). The process of dynamic rock pulverization (**Figures 2A–D**) consumes massive amounts of energy, e.g., in gouge formation it sums up to 50% of earthquake energy (Wilson et al., 2005). During grinding, most energy (~97%) is converted to heat, with only a small portion (<1%) actually contributing to fracturing (Spray, 1992).

#### Friction/Heat

Near the base of the moving rock mass, confining forces are largest and so the majority of frictional energy dissipation occurs in this zone (Pollet and Schneider, 2004). Disintegration and heating of the rock mass mainly arise (i) along well-defined persistent shear planes, or (ii) as a total disintegration of the whole mass. Shearing may be localized to a thin, discrete layer and frictional heating of bedrock may reduce basal strength (Hu et al., 2018, 2019; Hu and McSaveney, 2018). For (i), a high proportion of the energy release is focused on only a small proportion of the whole mass and will cause significant heating up to a partial melting of clasts, called *frictionite* (Heuberger et al., 1984; Erismann and Abele, 2001). As soon as particles are  $\sim 1 \ \mu m$  and below, the amount of heat produced by their elastic and plastic deformation leads to their melting (Spray, 2005). In rare cases (i) this heating can cause centimetre-thick melting of rock and formation of frictionites at temperatures of 1700°C (Erismann et al., 1977; Weidinger et al., 2014). Discrete layers of more intense fragmentation contain microbreccias and traces of partial melting (frictionite along shear planes; Schramm et al., 1998; Weidinger and Korup, 2009). For phase transitions, latent energy is absorbed. Besides, frictional shearing is controlled by the production and decay of random kinetic energy during gravitational work (Preuth et al., 2010). Random kinetic energy is referred to the random motion and inelastic interaction between the fragments; it is irreversible because it cannot perform mechanical work (Bartelt et al., 2006; Buser and Bartelt, 2009; Christen et al., 2010). For quantification, Schneider et al. (2010) argue that the total frictional work best correlates with the seismic signal of a rock (-ice) avalanche. The seismograph represents a small, but proportional fraction of this energy loss.

#### Erosion/Entrainment/Role of Water

There is an apparent increase in rock-avalanche mobility with volume (e.g., Heim, 1932; Scheidegger, 1973). The volume can be increased either by fragmentation up to 25-30% (Hungr and Evans, 2004; Crosta et al., 2007) or by the entrainment of substrate material. How enormous the effect of entrainment is, can be shown by the 2000 Tsing Shan event (Hong Kong), where a small volume of 150 m<sup>3</sup> of material grew to 1620 m<sup>3</sup> because of the strong erosion along the slope (Hungr and Evans, 2004).

Entrainment strongly depends on the character of the path material (Crosta et al., 2009; Aaron and McDougall, 2019) and, for example, may cause high basal shear resistance and momentum loss, when overrunning bedrock or dry bed material (Iverson et al., 2011; Aaron et al., 2017; Whittall et al., 2017; Aaron and McDougall, 2019). In other cases, basal friction is reduced and mobility enhanced (Hungr and Evans, 2004; Aaron and Hungr, 2016; Coe et al., 2016). On the one hand, entrainment is an energy sink because the erosion, uptake and incorporation of material along the travel path by plowing, scouring or even surficial scratching (Hu and McSaveney, 2018) is mechanical work, accompanied by heating. On the other hand, the gain in weight increases the energy budget by acting as an energy source and must not be neglected. Water plays an important role for the amount and rate of entrainment and erosion (Iverson and Ouyang, 2015). Especially for rock avalanches traveling on ice (Huggel et al., 2008; Deline et al., 2015; Bessette-Kirton et al., 2018; Walter et al., 2020) or wet, soft sediments (e.g., lake sediments, Figure 2E), the increased pore pressure enhances the scour of the bed, reduces basal friction and causes velocity, mass and momentum to increase (Iverson et al., 2011; Iverson, 2016; Johnson et al., 2016). Pure ice has a basal friction which is about 75% lower than that of pure rock. Hence, in rock-ice avalanches, a  $\sim 12.5\%$ reduction of basal friction angle is observed for every 10% increase in ice content (Sosio et al., 2012). The intergranular direct stress between single grains is reduced by pore-water pressure, i.e., in initially wet sediment more overburden is necessary to start fragmentation than in dry sediment (Abele, 1997). Water may escape quickly after deposition like at the Flims Rock Avalanche (Figure 2F; Pavoni, 1968; Calhoun and Clague, 2018) but increasing temperature may cause water to be pressurized (Voight and Faust, 1982) and/or vaporized, as it is proposed for the Vajont Rockslide (Habib, 1975) or the Köfels Rockslide (De Blasio and Medici, 2017). For the melting of ice, a specific latent heat of 334 kJ/kg is needed, and for vaporization  $\sim$ 2265 kJ/kg, which is almost seven times as much. There is a momentum exchange that consumes energy (Pudasaini and Krautblatter, 2014), and steam explosions are present, but probably camouflaged in the other energy dissipative processes. We have yet to understand the sedimentological imprint of steam explosions in the sediments.

## **Chemical Processes**

Chemical processes are often neglected in the energy balance of rock avalanches. Novel friction experiments on carbonate rocks, for example, show that at velocities of several meters per second carbon dioxide starts to degas due to thermal decomposition induced by flash heating after only a few hundred microns of slip (Mitchell et al., 2015). This process creates vesicular degassing rims in dolomite clasts and crystalline calcite cement (Anders et al., 2010; Mitchell et al., 2015) and may allow the upper rock mass to slide over a "cushion" of pressurized material. Around 800–850°C, talc and dolomite start to decompose (Hu et al., 2018) and to produce high-pressure live steam and carbon dioxide (Habib, 1975; Hu et al., 2019). De Blasio and Medici (2017) found bubbles grown in the frictionites of the Köfels Rockslide, which they ascribe to water vapor, either due to seeping of vadose water through rock fissures prior to the rock-slope failure, or due to dehydroxylation of the mica, which occurs at ~700°C (Alexiades and Jackson, 1967). Also, the existence and relative increase of pyrophyllite on sliding surfaces indicate hydrothermal alteration around 450°C (Schäbitz et al., 2015). The accumulation of pyrophyllite at the sliding surface results in reduced shear strength. Also, graphitization (crystallization of amorphous carbon) was recognized in slip zones as phase transformation, which implies frictional heating due to rapid sliding (Oohashi et al., 2011). Graphite is well known as an effective solid lubricant in fault zones with a friction coefficient as low as that of smectite,  $\mu = 0.1$  (Oohashi et al., 2014).

## **DISCUSSION OF RESEARCH NEEDS**

Processes during the disintegration phases of rock avalanches are beyond observation, and we have very few analogs that show pressure and temperature conditions inside rock avalanches. Thus, it is likely that we neglect important processes such as steam explosions, partial melting, chemical transitions, and material explosion processes at high pressures.

To systematically decipher relevant processes in rock avalanches, we propose that the energy balance needs to be considered more seriously, since it will help us to reveal energyrelevant processes that we would otherwise neglect. Here we propose to balance the primary energy input to the system constrained by the potential/kinetic energy of the moving rock masses ("Energy sources"). "Energy sinks" include heating, friction, inelastic collision, entrainment, compression during crustal deformation, chemical energy consumption, phase transition solid – fluid – gas, dust production and bouncing as well as sound and microseismicity, and momentum exchange between solid and fluid.

Using an energy balance approach, we can attribute proportions of the energy transmission to certain processes and we can rule out others. However, for this approach, we have to find ways to accurately constrain the 3D deposition temperature of the rock avalanche by new methods as has been exemplified in a few cases in this paper. The influence of the substrate on types and rates of energy dissipation during disintegration and during the flow represent major research gaps and ask for more studies. For the hazard assessment of rock avalanches, it makes sense not only to differentiate between energy sources and sinks, but also to separate processes that favor mobility and runout length from those which may consume or release energy but do not essentially contribute to the hazard potential. Furthermore, we need to transfer the achievements gained in qualitative assessment toward a more quantitative approach.

Future research in the field should focus on analyzing spatial patterns of disintegration using surface mapping and 3D subsurface reconnaissance of rock slide/avalanche deposits using geophysical methods at varying scales. Sedimentological analyses

reveal abundant information on internal processes, for instance high-stress comminution preserved in fine-sediment signatures (Reznichenko et al., 2012). There is a great demand for study cases with petrographic analysis at microscopic scale (Weidinger et al., 2014), and for such with cross sections through the debris (Locat et al., 2006).

Future research in the laboratory should focus on the implementation of disintegration scenarios in large-scale analog models to help better understand the impact of disintegration and heating on runout length. This way, a conceptual physical (and chemical) model of rock-avalanche disintegration in time and space may be set up in a first step, followed by the implementation in benchmark one- and two-phase runout models.

## CONCLUSION

- (1) Due to the law of energy conservation we have a superior tool to decipher processes we have yet neglected in rock avalanches: heating, friction, inelastic collision, entrainment, compression crustal deformation, chemical energy consumption, phase transition solid – fluid – gas, dust production and bouncing as well as sound and microseismicity generation and momentum exchange.
- (2) Energy dissipation is concentrated in the disintegration zone where energy estimations indicate considerable heating above 100°C of significant portions of the rock mass.
- (3) The spatial pattern of heating is characteristic for individual types of movement ranging from concentrated heating by friction along defined sliding planes to diffuse clustered heating in crushing zones near to obstacles.
- (4) Massive entrainment where large rockslides drive into, or override, valley sediments also evidently causes crushing and very likely significant heating.
- (5) Massive energy dissipation may leave a distinct sedimentological signal detectable in compressive and extensional flow structures, melting or new mineral formation, rock-avalanche structure, material composition, brecciation and fine-sediment signature.

## **AUTHOR CONTRIBUTIONS**

SK drafted the work, substantially contributed to the conception and design of the work, and acquisition and analysis of literature/data for the work. MK substantially contributed to the concept and revised the work critically for important intellectual content and provided approval for publication. Both authors contributed to the article and approved the submitted version.

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**Conflict of Interest:** The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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