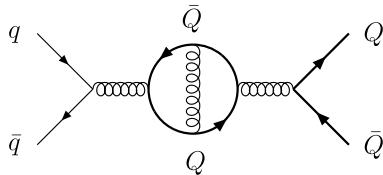




Corrigendum

Corrigendum to “Threshold expansion of the $gg(q\bar{q}) \rightarrow Q\bar{Q} + X$ cross section at $\mathcal{O}(\alpha_s^4)$ ” [Phys. Lett. B 690 (2010) 483–490]Martin Beneke^a, Michal Czakon^b, Pietro Falgari^c, Alexander Mitov^d, Christian Schwinn^b^a Physik Department T31, Technische Universität München, James-Franck-Straße 1, D-85748 Garching, Germany^b Institut für Theoretische Physik E, RWTH Aachen University, D-52056 Aachen, Germany^c Institute for Theoretical Physics and Spinoza Institute, Utrecht University, 3508 TD Utrecht, The Netherlands^d Cavendish Laboratory, JJ Thomson Avenue, Cambridge CB3 0HE, United Kingdom

Ref. [1] noted that there is an additional contribution to the singular terms in the threshold expansion of the heavy-quark pair production cross section at $\mathcal{O}(\alpha_s^4)$ from the one-particle reducible vacuum polarization contribution



to the 2-loop virtual correction, which was omitted in our calculation. This affects only the quark-antiquark production channel $q\bar{q} \rightarrow Q\bar{Q}$. In consequence the expression

$$\sigma_{q\bar{q},8}^{(2,0)} \Big|_{\text{VP}} = \left(C_F - \frac{C_A}{2} \right) 8\pi^2 \ln \beta = -13.160 \ln \beta$$

must be added to Eq. (5), which changes the coefficient of the $\ln \beta$ term in Eq. (8) from 528.557 to 515.397. Eq. (8) should read correctly

$$\begin{aligned} \sigma_{q\bar{q}}^{(2)} &= \frac{3.60774}{\beta^2} + \frac{1}{\beta} \left(-140.368 \ln^2 \beta + 32.106 \ln \beta + 3.95105 \right) \\ &\quad + 910.222 \ln^4 \beta - 1315.53 \ln^3 \beta + 592.292 \ln^2 \beta \\ &\quad + 515.397 \ln \beta + C_{q\bar{q}}^{(2)}. \end{aligned}$$

The new term amounts to a reduction of the coefficient of $\ln \beta$ by 2.5%. The numerical effect on the total cross section is negligible.

As discussed in Ref. [2] in the more general context of pair production of squarks and gluinos, the above additional term should be interpreted as a correction to the “Coulomb function” $J_{R_\alpha}(E)$ in Eq. (1) from an annihilation contribution, here from $Q\bar{Q} \rightarrow Q\bar{Q}$, to the NRQCD Lagrangian. Accordingly, the term

$$v_{\text{ann}}^{R_\alpha,S} (-D_{R_\alpha}) 8\pi^2 \ln \beta$$

should be added to the general result for $\sigma_X^{(2)}$ in Eq. (A.1). For a given heavy-particle state, $v_{\text{ann}}^{R_\alpha,S}$ depends on the colour representation R_α and total spin S of the pair. For $Q\bar{Q}$, only $v_{\text{ann}}^{8,1} = 2T_F = 1$ is different from zero. Results for $v_{\text{ann}}^{R_\alpha,S}$ as well as the quantity v_{spin} , which appears in Eq. (A.1), can be found in Ref. [2] for the case of pair production of squarks and gluinos.

References

- [1] P. Bärnreuther, M. Czakon, P. Fiedler, J. High Energy Phys. 1402 (2014) 078, arXiv: 1312.6279 [hep-ph].
- [2] M. Beneke, J. Piclum, C. Schwinn, C. Wever, J. High Energy Phys. 1610 (2016) 054, arXiv:1607.07574 [hep-ph].