



## High performance Integrated Electro-Hydraulic Actuator for robotics. Part II: Theoretical modelling, simulation, control & comparison with real measurements<sup>\*</sup>

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### abstract

The paper deals with the theoretical modelling of the high performance Integrated Electro-Hydraulic Actuator (IEHA) introduced in the first part of this work. Based on the hydraulic and dynamic equations, we carried out an analysis of the IEHA input/output relation. A simulation based on this input/output model is detailed showing the ability to access several time derivatives of the payload position mainly its “jerk”, which represents an important issue whenever motion control is considered. We show the theoretical dynamic properties of the proposed solution through open and closed loop simulation results based on Matlab-Simulink. A comparison between the theoretical and experimental results obtained by a built prototype is given in order to demonstrate the promising performances of the proposed solution.

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### 1. Introduction

In the first part of this work [1], a novel Integrated Electro-Hydraulic Actuator (IEHA) was proposed. The IEHA is basically based on a technology named Hydrostatic Transmission, first proposed for robotic application by Bobrow et al. [2] and followed by Habibi who introduced the concept of ElectroHydraulic Actuator (EHA) [3]. In all hydraulic solutions for robotic applications [4–7], the hydraulic motor control and hence the payload or the positioning is carried out by controlling either a fixed or variable displacement pump rotation. In the case of a fixed displacement pump, a power control of the servo electric motor has to be carried out, leading to sophisticated control laws which limit the use of this kind of solution in high numbers system, such as humanoid robots. In the case of a variable displacement pump with a constant speed electric motor, a bidirectional pump is also needed leading to the necessity of inverting the electric motor direction, which induces the appearance of a dead-band, thus, a loss in the system's effi-

ciency. Consequently, in all former solutions, power production is made instantaneously and the hydrostatic actuator has to be sized according to the worst-case constraints. Our approach in developing the IEHA was inspired by humans' behavior in their process of motion/force production. It is well known that the human being can achieve with their muscles two kinds of modes called isometric and isotonic [8]. For the isometric functioning, the link (arm or leg) has to keep a given position and to exert some force varying from zero (without contact with the environment) to a maximum value (contact force, pushing, pulling or propelling during locomotion gait). The force variation is ensured by the recruitment of muscular fibers with some electric activation signal. For the isotonic mode, the muscles are compensating the gravity effects and no more contraction (activation) is required to carry out smooth movement. In the IEHA proposed solution, a built-in micro-valve activated with a magnetic coil and controlled with a Linear Variable Differential Transformer (LVDT) sensor [9] is integrated in order to allow us to emulate the muscle's activation phenomena by changing the eccentricity of the micro-pump. Inverting the flow and the pressure direction by keeping the same rotation is carried out with a built-in rotary passive distributor, which has three possible positions depending on the pressure in the input and output lines of the micro-pump. The established mathematical model of the input/output relation shows the capacity of the system to access to the “jerk” of the payload and giving us the possibility to minimise it whenever carrying isotonic movements.

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### Nomenclature

E	eccentricity (m)
$S_p$	piston surface ( $m^2$ )
Q	micro-pump average flow ( $m^3/min$ )
$N_p$	global pistons number (–)
$\omega$	rotation speed (rpm)
P	fluid pressure (bar)
X	position of the micro-valve body (m)
Y	payload position (m)
R	theoretical micro-valve radius (m)
$\sigma$	micro-valve surface opening ( $m^2$ )
$V_{P_b}$	fluid velocity at point $P_b$ ( $m/s^2$ )
$V_{D_1}$	fluid velocity at point $D_1$ ( $m/s^2$ )
	fluid volumic density ( $kg/m^3$ )
$P_0$	micro-valve fluid source pressure (bar)
$\phi$	difference of the pressure created by the movement of the micro-valve (bar)
$Q_e$	flow created by the movement of the micro-valve ( $m^3/s$ )
C	Vena Contracta constant (–)
$V_{e1}$	$CH_{e1}$ chamber's volume ( $m^3$ )
	compressibility coefficient ( $bar^{-1}$ )
$S_v$	active section of the linear actuator ( $m^2$ )
$V_v$	$CH_v$ chamber's volume ( $m^3$ )
$F_{e1}, F_{e2}$	forces produced by the fluid of the micro-valve on the barrel (N)
$F_{pi}$	force exerted by the $i$ th piston at the barrel (N)
$P_{v1}$	$CH_v$ chamber's pressure (bar)
$P_{pi}$	pressure of the $i$ th piston (bar)
$m_e$	barrel mass (kg)
$\ddot{x}_3$	barrel acceleration ( $m/s^2$ )
$S_e$	barrel active surface ( $m^2$ )
$\omega$	rotation speed (rad/s)
$V_{pi}$	fluid velocity in chamber $CH_{pi}$ (m/s)
$L_0$	linear actuator total length (m)
m	payload mass (kg)
$F_c$	external force applied to the payload (N)

This paper is organised as follows: in Section 2, the proposed IEHA input–output relationship is developed. Based on this theoretical model, a simulation tool using Matlab-Simulink is carried out in the third section. Finally, comparisons between theoretical and experimental results are detailed and further developments of this work are given.

## 2. IEHA input–output relationship

In this section, the relationship of IEHA input–output is established. This relationship is linked to the input, the position  $X$  of the micro-valve body to the payload position  $Y$  in the case of a linear actuator. The same relationship can be established even if a rotary hydraulic actuator is used. In this case, the input–output is linked to  $X$  with the angular position ( $\theta$ ) of the actuator. Fig. 1 shows a simplified sketch of a planar section of the proposed IEHA, as in the case of a linear actuator. Based on these relationships, models of the different stages of the IEHA can be established. We start the modelling process with the micro-valve stage, as it controls the input signal  $X$ . Then, the equation of motion of the micro-pump barrel is established using the average flow calculation given by the following equation:

$$Q = 2N_p E S_p Q \quad (1)$$

where  $N_p$  the total number of pistons,  $E$  is the eccentricity of the micro-pump,  $S_p$  is the surface of one piston and  $\omega$  is the rotation speed of the input shaft. A microscopic study to quantify the amount of each piston contribution can be carried out, in order to provide a better approximation of the dynamic effects applied to the micro-pump barrel. Lastly, we model the final stage of the process, which deals with the dynamics of the linear actuator driving the payload.

### 2.1. Micro-valve equation of motion

The opening,  $\sigma$ , allows the fluid passage is defined by the surface of the cylindrical opening created by the movement of the body  $C_4$  (see Fig. 1). If there is a symmetry of manufacturing in the micro-valve, the movement of the body  $C_4$  towards the left or towards the right yields the same  $\sigma$  at both points  $O_1$  and  $O_2$ . Theoretically, the micro-valve body is assumed to be a perfect cylinder of radius  $R$  leading to the following expression of  $\sigma$ :

$$\sigma = 2nR(X - E) \quad (2)$$

As we supposed the friction is negligible, if the pressure  $P_d$  is assumed to be equal to  $P_0$  and the pressure  $T_d$  is considered to be equal to the atmospheric pressure, then the pressure in chambers  $CH_{e1}$  and  $CH_{e2}$  is thus,  $\frac{P_0}{2}$  in the rest position. According to the Bernoulli theorem, for a given point  $P_b$  of a fluid, it is possible to write the following relation:

$$\frac{V_{P_b}^2}{2} + \frac{P}{\rho} + gh = C^{te} \quad (3)$$

where  $V_{P_b}$  is the fluid velocity at point  $P_b$ ,  $P$  is its pressure at this point,  $\rho$  represents the fluid volumic density,  $h$  is the height of the point  $P_b$  with regard to a given reference and  $g$  the gravity acceleration. By applying this theorem between both points  $O_1$  and  $D_1$  while assuming that these points are for the same height, the velocity of the fluid  $V_{D_1}$  at the point  $D_1$  is give by:

$$V_{D_1} = \frac{\sigma \sqrt{P}}{2} \quad (4)$$

As the micro-valve body begins to move towards the right, as the pressures in chambers  $CH_{e1}$  and  $CH_{e2}$  are no longer equal. The pressure is then equal to  $((P_0/2) + \phi)$  in the  $CH_{e1}$  chamber and to  $((P_0/2) - \phi)$  in the  $CH_{e2}$  chamber.  $\phi$  represents the difference of the pressure created by the movement of the micro-valve body. Hence, the pressure difference  $\Delta P$  between both points  $O_1$  and  $D_1$  is given by:  $\Delta P = P_0 - \frac{P_0}{2} + \phi = \frac{P_0}{2} - \phi$ . Moreover, using Vena Contracta constant,  $C$ , which is almost equal to 0.62, the flow  $Q_e$  created by the micro-valve can be defined by:

$$Q_e = C \cdot V_{D_1} \quad (5)$$

By substituting Eqs. (4) and (2) in Eq. (5), the flow  $Q_e$  can be established by:

$$Q_e = K_1(X - E) \frac{2((P_0/2) - \phi)}{\rho} \quad (6)$$

with  $K_1 = 2nC R$ . The flow  $Q_e$  can also be established by using its definition and the fluid compressibility with a second relation as:

$$Q_e = S_e \dot{E} + V_e \dot{\phi} \quad (7)$$

where  $\dot{E}$  is the time derivative of the eccentricity  $E$ ,  $V_e$  represents the  $CH_{e1}$  chamber's volume and  $\beta$  is the compressibility coefficient defined as:

$$\beta = -\frac{1}{V} \frac{\partial V}{\partial P} \quad (8)$$

$V$  and  $P$  are, respectively, the volume and the pressure. For mineral fluid, this coefficient is almost equal to  $1/16,000 \text{ bar}^{-1}$ .

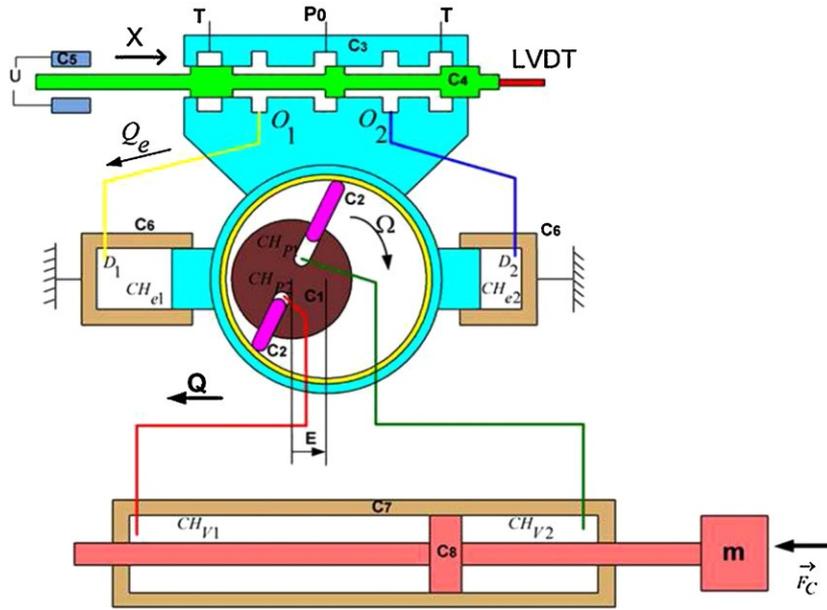


Fig. 1. IEHA simplified model (a linear actuator).

2.2. Micro-pump barrel dynamic equilibrium

The average flow value given by Eq. (1) can be also expressed by using the relation between the actuator linear output velocity  $\dot{Y}$  and coefficient as:

$$Q = S_v \dot{Y} + V_v P_{v1} \quad (9)$$

where  $S_v$  is the active section of the linear actuator,  $V_v$  is the  $CH_v$  chamber's volume and  $P_{v1}$  is the pressure in this chamber (see Fig. 1). Applying Newton's equation to the micro-pump barrel ( $C_3$ ), with mass is  $m_e$  and acceleration of  $\ddot{x}_3$ , all external forces can be identified. These forces, as depicted in Fig. 2, can be determined by the following:

1. Two forces produced by the fluid of the micro-valve stage on the barrel are  $F_{e1}$  and  $F_{e2}$ .
2. Force exerted by the  $i$  th piston and induced by pressure  $P_{pi}$  in the corresponding chamber  $CH_{pi}$ . This force is given by  $F_{pi}$  and has a magnitude of  $S_p P_{pi}$ .
3. Gravity force which is perpendicular to the micro-pump shaft.

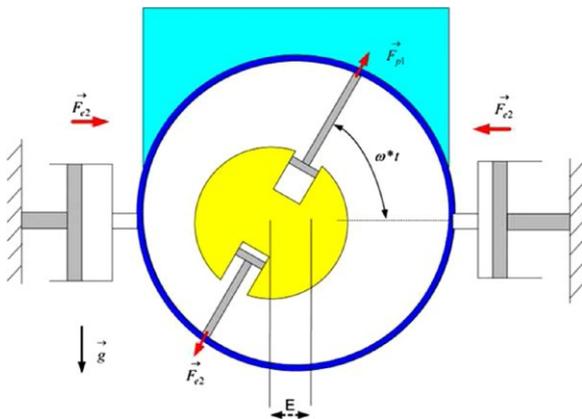


Fig. 2. Applied forces on the pump-barrel.

Hence, Newton's equation in its vectorial form can be expressed as follows:

$$m_e \ddot{x}_3 = F_{e1} + F_{e2} + F_{p1} + F_{p2} + m_e g \quad (10)$$

Using the pressures in both chambers  $CH_{e1}$  and  $CH_{e2}$ , the scalar projection of the above equation on the shaft axis leads to the following formulation:

$$m_e \ddot{E} = S_p P_{p1} \cos(\omega t) + S_p P_{p2} \cos(\omega t + n) + (P_0 + \phi) S_e - (P_0 - \phi) S_e \quad (11)$$

The later equation can take the following form:

$$m_e \ddot{E} = S_p (P_{p1} - P_{p2}) \cos(\omega t) + 2\phi S_e \quad (12)$$

2.3. Linear actuator dynamics

Applying Bernoulli theorem between the pistons located on the side of the micro-pump and the actuator chambers, for both cases, aspiration and expulsion leads to the following set of equations:

$$\frac{V_{p1}^2}{2} + \frac{P_{p1}}{\rho} = \frac{\dot{Y}^2}{2} + \frac{P_{v1}}{\rho} \quad (13)$$

$$\frac{V_{p2}^2}{2} + \frac{P_{p2}}{\rho} = \frac{\dot{Y}^2}{2} + \frac{P_{v2}}{\rho} \quad (14)$$

where  $V_{pi}$  is the fluid velocity in chamber  $CH_{pi}$ , which can be assumed to be equal to the linear velocity of the corresponding piston  $pi$ . Let us consider a simplified system with two pistons  $p_1$  and  $p_2$  (one in aspiration mode while the other in expulsion mode) and located, respectively, at  $(l_1 = \omega t)$  and  $(l_2 = \omega t + n)$ . Hence, and for a given eccentricity  $E = c^{le}$ , the linear velocities  $V_{p1}$  and  $V_{p2}$  can be established by the time derivative of the pistons position  $L_p$  as follows<sup>1</sup>:

$$V_{p1} = +E\omega \sin(\omega t) - \frac{E^2 \omega \sin(\omega t) \cos(\omega t)}{R_b^2 - E^2 \sin^2(\omega t)} \quad (15)$$

<sup>1</sup> Is clear that same results can be derived for each layer  $n_e$ , with odd number  $n_p^{max}$  of pistons for each.

$$V_{p2} = -E\omega \sin(\omega t) - \frac{E^2\omega \sin(\omega t) \cos(\omega t)}{R_b^2 - E^2 \sin(\omega t)^2} \quad (16)$$

Going back to both Eqs. (13) and (14) and subtracting one from the other leads to:

$$\tilde{\mathcal{P}}_v = \tilde{\mathcal{P}}_p + \frac{1}{2}(V_{p2}^2 - V_{p1}^2) \quad (17)$$

This equation can be rewritten using both Eqs. (15) and (16), which gives the following non-linear equation:

$$\tilde{\mathcal{P}}_v = \tilde{\mathcal{P}}_p + 4 \frac{E^3\omega \sin(\omega t)^2 \cos(\omega t)}{R_b^2 - E^2 \sin(\omega t)^2} \quad (18)$$

The second term on the right side produces non-desired non-linearity, therefore, based on comparison between the various geometrical sizes, we established the following statement:

$$\frac{E}{R_b} < 1 \Rightarrow E\omega \sin(\omega t) \gg \frac{E^2\omega \sin(\omega t) \cos(\omega t)}{R_b^2 - E^2 \sin(\omega t)^2} \quad (19)$$

It is then possible to assert that  $V_{p2} \approx -V_{p1}$ , which leads us to establish the final form of the Bernoulli theorem for the pistons and the output actuator as follows:

$$\tilde{\mathcal{P}}_v = \tilde{\mathcal{P}}_p \quad (20)$$

Once again, applying Newton's equation to the payload, which its mass is equal to  $m$  while taking into account the external applied force  $F_c$ , leads to the following relationship:

$$m\ddot{Y} = S_v \tilde{\mathcal{P}}_v - F_c \quad (21)$$

#### 2.4. Input–output relation

Substituting Eqs. (20) and (21) into Eq. (12) gives the following formulation:

$$\frac{S_p}{S_v}(m\ddot{Y} + F_c) \cos(\omega t) + 2\phi S_e = m_e \ddot{E} \quad (22)$$

Since our objective is to establish the relation between the input  $X$  and the output  $Y$ , it is necessary to get the following expression of  $\phi$  from the above Eq. (22):

$$\phi = K_2 \ddot{E} - K_3 \ddot{Y} \cos(\omega t) - K_4 F_c \cos(\omega t) \quad (23)$$

where constants  $K_i, i=2, 3, 4$  are introduced in order to simplify the notations:

$$K_2 = \frac{m_e}{2S_e}, \quad K_3 = m \frac{S_p}{2S_e S_v}, \quad K_4 = \frac{S_p}{2S_e S_v} \quad (24)$$

$$\begin{aligned} K_5 \ddot{Y} + mK_6 Y^{(4)} + K_6 \ddot{F}_c + mK_7 \dot{Y} Y^{(3)} + mK_7 Y Y^{(4)} + K_7 \dot{Y} \ddot{F}_c + K_7 Y \ddot{F}_c &= \frac{K_1}{S_e} (X - K_5 \dot{Y} + mK_6 Y^{(3)} + K_6 \ddot{F}_c + mK_7 Y Y^{(3)} + K_7 Y \ddot{F}_c) \\ \times \left[ \frac{P_0}{2} - K_2 K_5 Y^{(3)} - mK_2 K_6 Y^{(5)} - K_2 K_6 F_c^{(3)} - mK_2 K_7 \ddot{Y} Y^{(3)} - mK_2 K_7 Y Y^{(4)} - mK_2 K_7 \dot{Y} Y^{(4)} - mK_2 K_7 Y Y^{(5)} - K_2 K_7 \ddot{Y} \ddot{F}_c \right. \\ \left. - K_2 K_7 \dot{Y} \ddot{F}_c - K_2 K_7 Y \ddot{F}_c - K_2 K_7 Y F_c^{(3)} + K_3 \ddot{Y} \cos(\omega t) + K_4 \ddot{F}_c \cos(\omega t) \right] \end{aligned} \quad (33)$$

When the micro-pump stage starts moving the output actuator from its equilibrium position, the pressure in both chambers  $CH_{v1}$  and  $CH_{v2}$  will change symmetrically allowing us to admit that  $P_{v1} = -P_{v2}$  or even more  $\tilde{\mathcal{P}}_v = 2P_{v1}$ . Hence, using Eqs. (1) and (9), the eccentricity  $E$  can be expressed as follows:

$$E = \frac{S_v}{2N_p S_p Q} \dot{Y} + \frac{V_v}{2N_p S_p Q} P_{v1} \quad (25)$$

The time derivative of Eq. (21) combined with the expression of  $\tilde{\mathcal{P}}_v$  leads to:

$$\dot{P}_{v1} = \frac{m}{2S_v} Y^{(3)} + \frac{\dot{F}_c}{2S_v} \quad (26)$$

$Y^{(k)}$  is the  $k$ th time derivative of  $Y$  (i.e.  $Y^{(3)}$  is the pay-load jerk). Let us assume that the output actuator initial position is in the middle of its stroke and its total length is  $L_0$ . In that case, the volume of the chamber is given by:

$$V_v = \frac{L_0}{2} + Y S_v \quad (27)$$

This leads us to formulate Eq. (25) as:

$$E = K_5 \dot{Y} + mK_6 Y^{(3)} + K_6 \dot{F}_c + mK_7 Y Y^{(3)} + K_7 Y \dot{F}_c \quad (28)$$

where constants  $K_i, i=5, 6, 7$  are introduced for the above mentioned reason:

$$K_5 = \frac{S_v}{2N_p S_p Q}, \quad K_6 = \frac{L_0}{8N_p S_p Q}, \quad K_7 = \frac{1}{4N_p S_p Q} \quad (29)$$

If we consider that the micro-pump shaft is rotating at almost constant speed ( $\omega \diamond C^t$ ), then the following relationship can be written:

$$\dot{E} = K_5 \ddot{Y} + mK_6 Y^{(4)} + K_6 \ddot{F}_c + mK_7 \dot{Y} Y^{(3)} + mK_7 Y Y^{(4)} + K_7 \dot{Y} \ddot{F}_c + K_7 Y \ddot{F}_c \quad (30)$$

Taking the time derivative of the above equation leads to the expression of the micro-pump barrel's acceleration  $\ddot{E}$ , which can be established as follows:

$$\begin{aligned} \ddot{E} = K_5 Y^{(3)} + mK_6 Y^{(5)} + K_6 F_c^{(3)} + mK_7 \ddot{Y} Y^{(3)} + mK_7 \dot{Y} Y^{(4)} \\ + mK_7 Y Y^{(4)} + mK_7 Y Y^{(5)} + K_7 \ddot{Y} \ddot{F}_c + K_7 Y \ddot{F}_c + K_7 Y \ddot{F}_c + K_7 Y F_c^{(3)} \end{aligned} \quad (31)$$

On the other hand, the leakage in the chambers ( $CH_{ei}, i=1, 2$ ) can be neglected since their volume is smaller than of the output actuator chambers. By substituting Eqs. (23) and (7) in Eq. (6), the later can be rewritten in the following form:

$$S_e E = K_1 (X - E) \frac{P_0}{2} - K_2 \ddot{E} + K_3 \ddot{Y} \cos(\omega t) + K_4 F_c \cos(\omega t) \quad (32)$$

A non-linear relationship links an implicit form of the input  $X$  to the output  $Y$  through the times derivative of  $Y$  and the external force  $F_c$  applied to the pay-load. To get the full expression, one can replace the above equation by the internal variable  $E$  (Eq. 28) and its first and second time derivatives  $\dot{E}$  and  $\ddot{E}$  given, respectively, by Eqs. (30) and (31). This leads to a complex relation, which is the base of the simulation carried out in the following section.

### 3. Simulation results

Given the complexity of the proposed system, to achieve a globally optimised solution it is important to proceed in an iterative manner. This iterative process is coupled with analyses by simulations of a virtual model and experimental measurements taken with a built prototype (described in the next section). The IEHA theoretical performances are established taking advantage of the

simulation results based on the relationships developed in previous section. The aim of the carried out simulations is to allow us to quantify the contributions of each IEHA stages (micro-pump, micro-valve and passive distributor). The forward objective of this simulation is double in the sense that it consists of comparisons between the simulated performances and the real measurements. Providing suggestions for improvements of the physical prototype, as well as, refinements of the several models included in the simulator. To this aim, two simulation models under Matlab-Simulink were developed. The relatively simplistic one has the objective of reproducing the functioning of the IEHA open-loop. The second aims at quantifying the IEHA behavior when closed loop control is carried out with sensors signals.

3.1. Open-loop model

As discussed in the previous section, the IEHA input–output relationship is a non-linear fifth order differential equation linked to the micro-valve position  $X$  and to the linear actuator position  $Y$ . The first simulation model is built by reducing the study to the simple case without taking into account the compressibility of the fluid. Fig. 3 presents a block diagram of the implemented simulator. In this model, the micro-valve block is built using two inputs  $X$  and  $\psi$  and an output which is the eccentricity  $E$ . Eq. (6) giving the flow  $Q_e$  in the case of neglecting the fluid compressibility is implemented in this block.

$$Q_e = 2 \cdot CR(X - E) \sqrt{\frac{2((P_0/2) - \psi)}{\rho}} \quad (34)$$

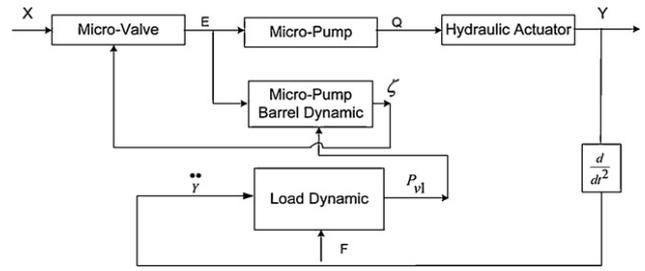


Fig. 3. IEHA simplified simulation model.

The micro-pump block presents the stage of hydraulic energy production implemented using Eq. (1), providing the average flow  $Q$ . Its input is the eccentricity  $E$  while its output is the above mentioned flow  $Q$ . By neglecting the fluid compressibility, the block simulate a linear actuator with the payload based on the following new Eq. (9). This simulation block has as input the flow  $Q$  and as output  $Y$ .

$$Q = S_v \dot{Y} \quad (35)$$

The dynamics of micro-pump barrel is implemented according to Eq. (12), in the simulation block with two inputs: the main one which is the eccentricity  $E$  and an auxiliary input providing the pressure of the first chamber of the linear actuator  $P_{v1}$ . The output of this block is the pressure of the chamber of the micro-pump barrel  $\psi$ . The auxiliary input is taken from the linear actuator’s chamber, as the pressure loss between the micro-pump and the linear actuator can be neglected (see Eq. (20)). As the passive distributor always guarantees that the connection of the return of the linear actuator

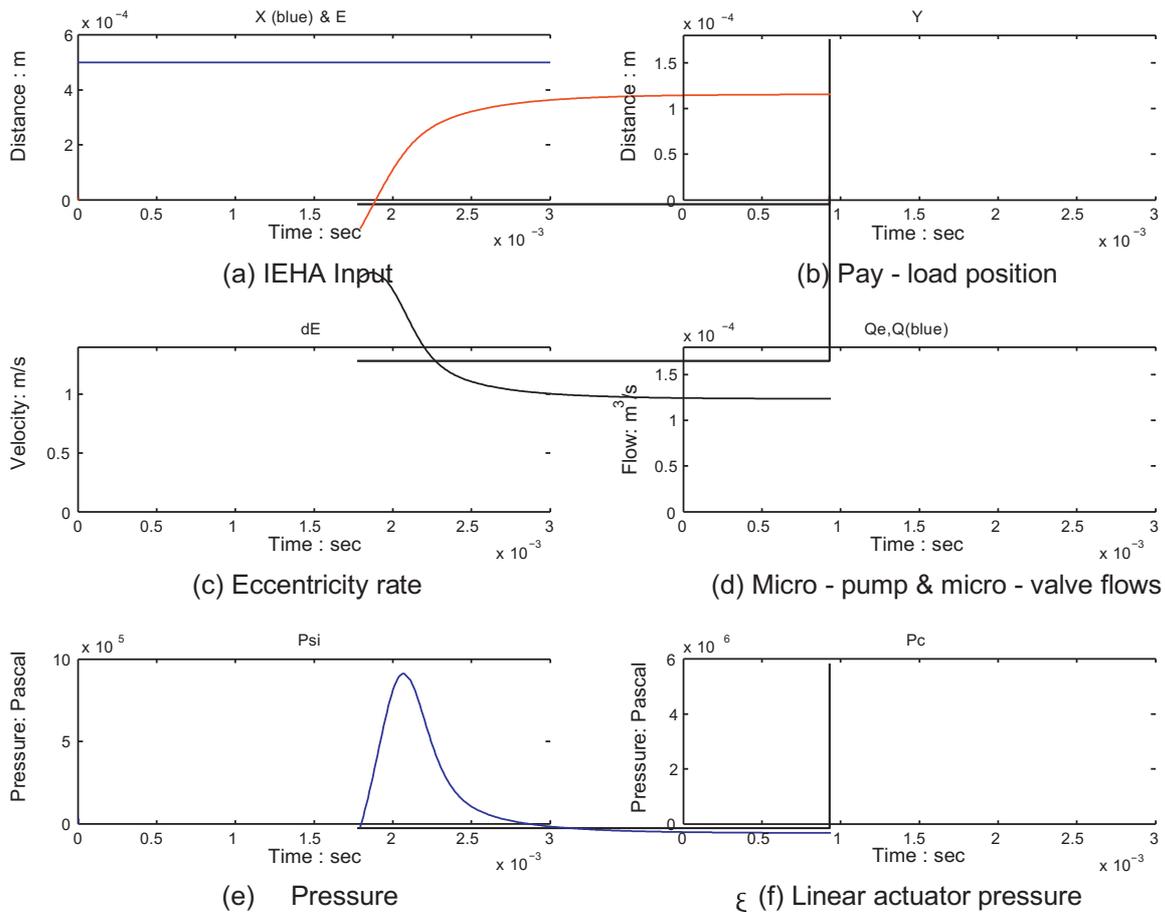


Fig. 4. Open-loop simulation results for an input step  $X$ .

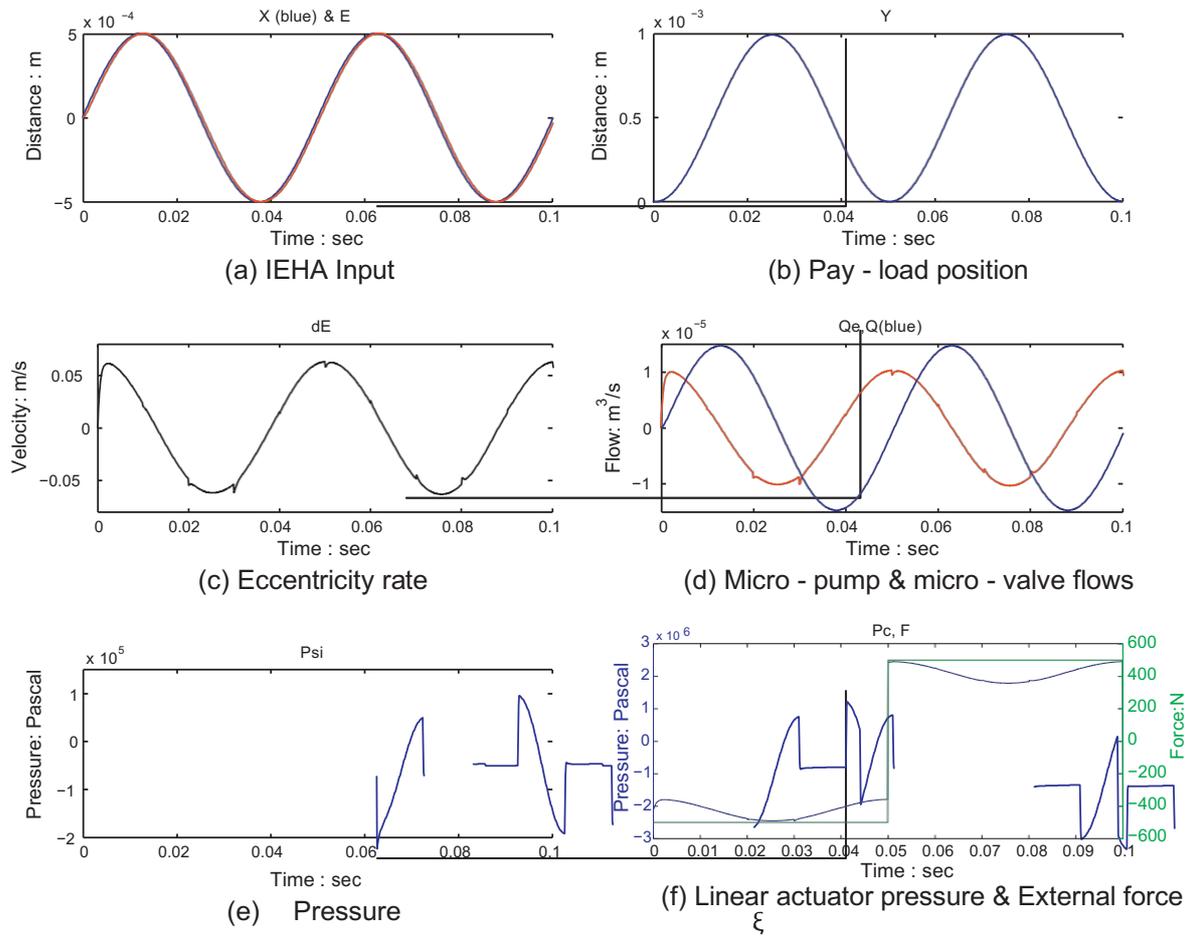


Fig. 5. Open-loop simulation results for external applied force.

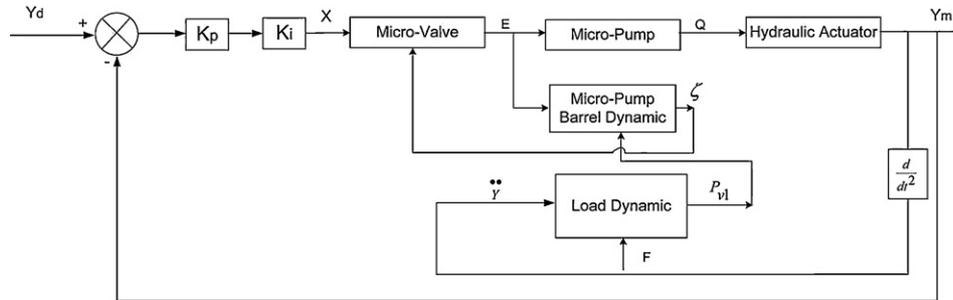


Fig. 6. IEHA closed loop simulation model.

$P_{v2}$  is equal to the atmospheric pressure. Hence, the pressure  $P_{v2}$  can be neglected, allowing us to define the considered block using the following equation:

$$m_e \ddot{E} = S_p P_{v1} \cos(\omega t) + 2 \zeta S_e \quad (36)$$

The last simulated block is concerned with the dynamics of the payload  $m$ , it has two main inputs  $\ddot{Y}$  and  $F_c$  and a single output, which is the pressure  $P_{v1}$ . It is implemented using Eq. (37):

$$m \ddot{Y} = S_v P_{v1} - F_c \quad (37)$$

Once all the simulation blocks are defined, the links between inputs and outputs are realised as shown in Fig. 3. Table 1 provides the values of the parameters used in the developed simulation model.

Table 1  
Parameters values used in the open-loop simulation.

Parameter	Physical meaning	Value
$m$	Payload mass	10 kg
$S_v$	Linear actuator active surface	2.35 cm <sup>2</sup>
$F$	External force applied to the payload	0 N
$m_e$	Micro-pump barrel mass	0.09 kg
$n_p$	Number of pistons	15
$S_p$	Surface of the piston	0.19 cm <sup>2</sup>
$S_e$	Micro-pump barrel active surface	1.64 cm <sup>2</sup>
$E_{max}$	Maximal eccentricity	0.05 cm
$r_{rig}$	Micro-valve radius	0.25 cm
$\omega$	Shaft speed	3000 rpm
$C$	Vena Contracta coefficient	0.62
	Fluid density	840 kg/m <sup>3</sup>
$P_0$	Micro-valve input pressure	10 <sup>5</sup> Pa

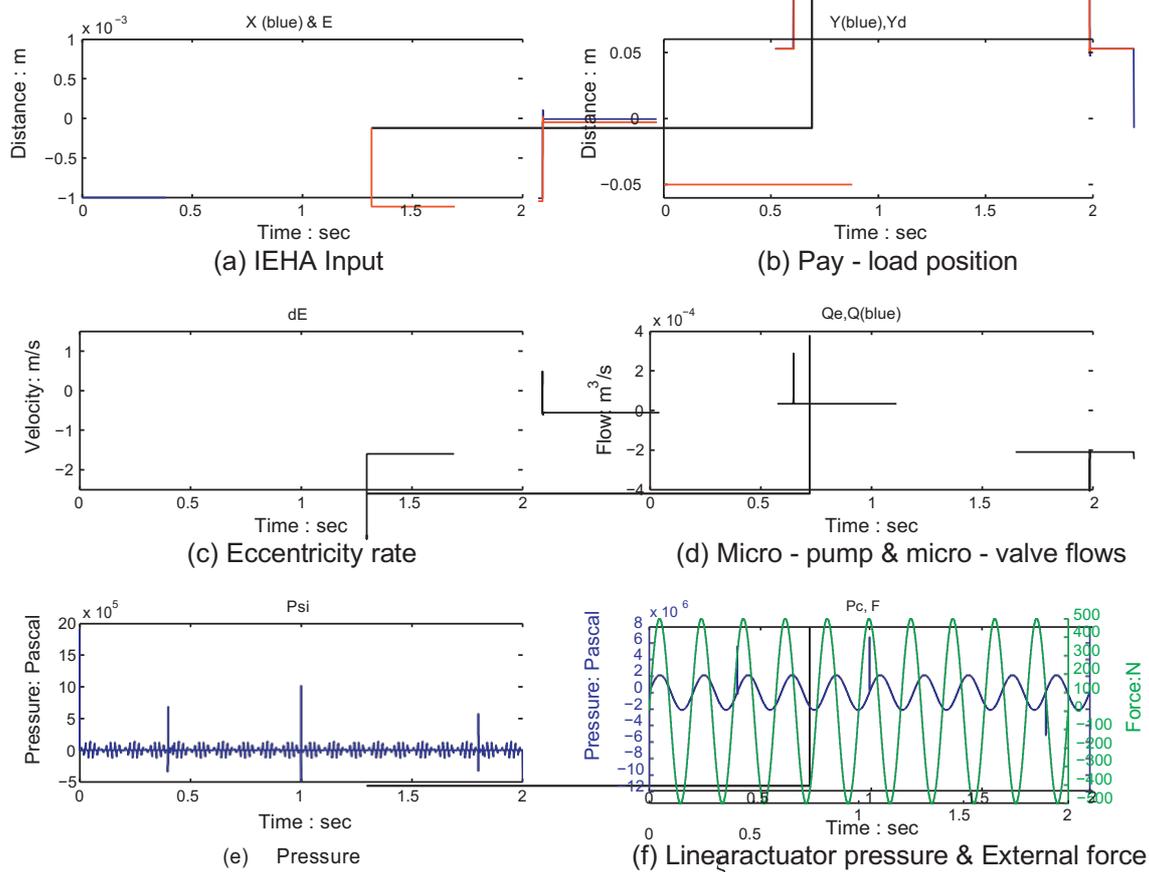


Fig. 7. Closed loop simulation results with external force and  $X_{\max} = 1$  mm.

### 3.2. Open-loop simulation results

Fig. 4 provide the results obtained from open-loop simulations. The input signal is at steps of 0.5 mm corresponding to the variable  $X$  controlling the position of the micro-valve. Since it is an open-loop simulation, the output  $Y$  increases infinitely and this for an eccentricity fixed to a maximal value. The time response of the micro-valve is almost of 1.5 ms. In this first simulation, no force is applied to the payload. This explains that the pressure converges to zero, as the payload's acceleration nullified.

The second simulation is carried out by applying a sinusoidal input with a period of 50 ms and of an amplitude of 0.5 mm. A periodic external force is also applied to the payload. This force has a period of 100 ms and an amplitude of 40 kg. Fig. 5 shows the resulting response of the system, which demonstrates clearly that the proposed IEHA has two intrinsic feedback loops. The first loop exists between the input  $X$  and the eccentricity  $E$ , while the second links the external applied force  $F_c$  to the pressure produced by the micro-pump  $P_{v1}$ .

### 3.3. Closed loop control

To be able to control the linear actuator output  $Y$ , a LVDT displacement sensor is used to measure the micro-valve position. A closed loop with a Proportional and an Integral terms are incorporated into the simulation model developed. The new simulation model is shown in Fig. 6. The values of the proportional and integral gains ( $k_p$ ,  $k_i$ ) are adjusted to have an acceptable time response. The values obtained after several trials are  $k_p = 0.01$  and  $k_i = 0.3$  (other parameters are given in Table 1). Fig. 7 shows that the system's time response is almost equal to 400 ms. By analysing the variation of the input  $X$  and of the eccentricity  $E$  over time, a very low

latency between these two variables appears. Nevertheless, both variables remain for a long time equal to their maximum values. This is probably due to the fact that the simulated micro-valve is not large enough.

By using a micro-valve with a bigger stroke ( $X_{\max} = 1.5$  mm), the time response is reduced to 100 ms, with the same values of the above mentioned gains ( $k_p = 0.01$  and  $k_i = 0.3$ ). Fig. 8 shows the new results. An alternative way to improve the time response of the system consists of increasing the micro-valve diameter which would increase its section, and hence, its dynamics.

## 4. Comparisons with experimental results

In order to carry out experimental results, the test-bed device dedicated to the IEHA [1] is used. This test-bed allows us to carry out continuously several experiments as detailed in the following sections:

### 4.1. Experiment 1: flow evolution with the rotation speed

Fig. 9 shows the flow variation produced by the IEHA according to the rotation speed of the input electric motor. Two curves are drawn on this figure. The first one (in black) concerns the theoretical flow calculated by means of the Eq. (1); the second (in magenta) is relative to the real flow measured by the test-bed (For interpretation of the references to color in this figure, the reader is referred to the web version of the article.). The difference between the experimental and theoretical results is greatest at low rotation speeds. This difference tends to decrease when the rotation speed of the input electric motor increases. This phenomenon is caused by the decrease of the centrifugal forces when the rotation speed decreases. The result demonstrated in this experiment

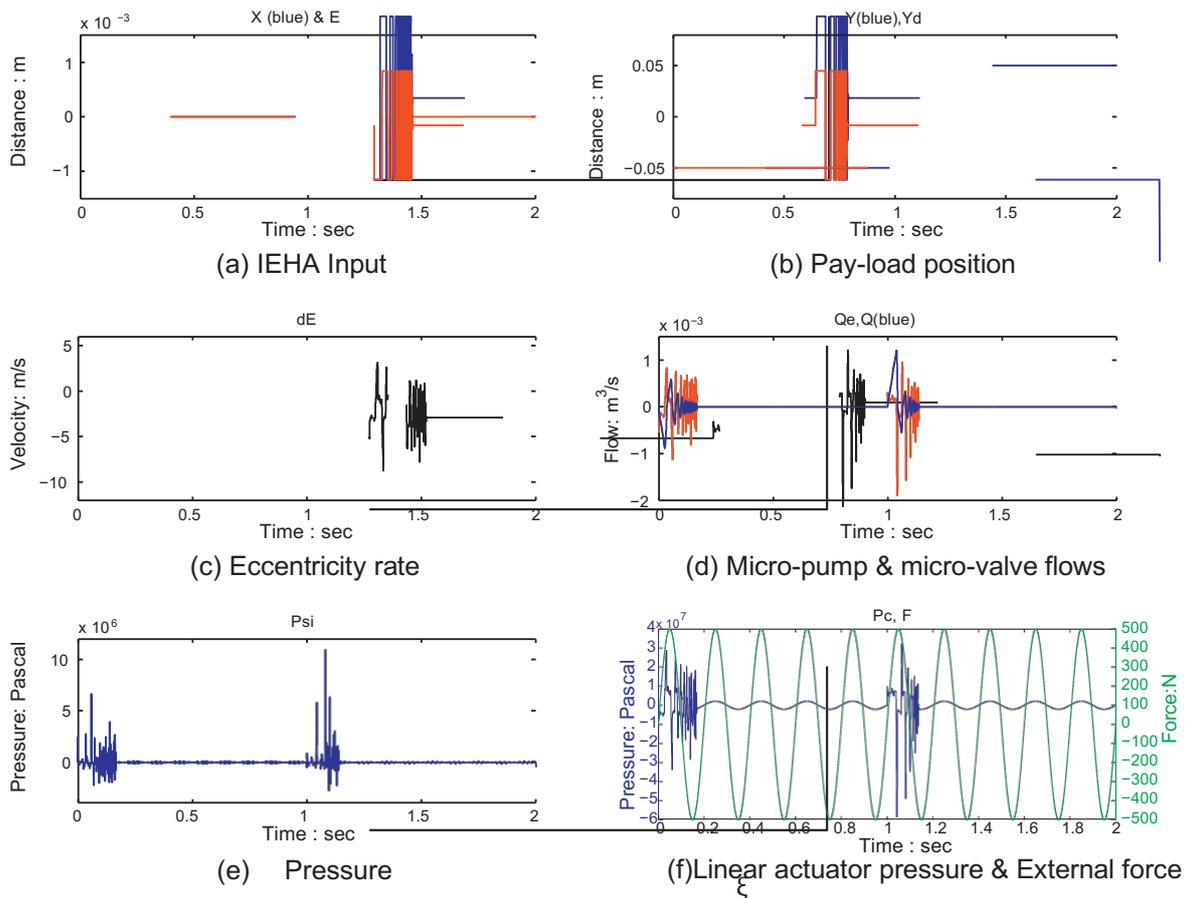


Fig. 8. Closed loop simulation results with external force and  $X_{max} = 1.5$  mm.

shows that there is an “optimal” rotation speed for which the centrifugal forces applied to pistons are big enough to allow these to accomplish their over-all displacement, and hence, to make the real flow rate to approach the theoretical value. Going over this optimal rotation speed, the increase of the flow is very low regards to the drawbacks related to the noise generated by the high speed. Practically, the optimal rotation speed is identified for this prototype around 3000 rpm.

4.2. Experiment 2: relation between the output pressure and the produced flow

In this experiment, the rotation speed  $\omega$  is fixed at 3000 rpm. The eccentricity  $E$  is equal to 0.13 mm and the number of used pistons is equal to 15. Every piston has a radius of in 0.25 mm. By changing the position of the flow reducer, the IEHA feels its payload increasing with an increase of its output pressure. Fig. 10 gives the flow variation produced for several values of output pressure (i.e. the payload), the two curves shown are: a theoretical one (in black) and the experimental one (in magenta) (For interpretation of the references to color in this figure, the reader is referred to the web version of the article.). The results show that when the output pressure increases, the produced flow decreases, and the distance between the real and theoretical values increases. Since all the characteristics of the actuator ( $E$ ,  $N_p$  and  $\omega$ ) are fixed, the flow lost takes place inside the IEHA due the internal leakages. This issue will be improved in the future planned prototypes.

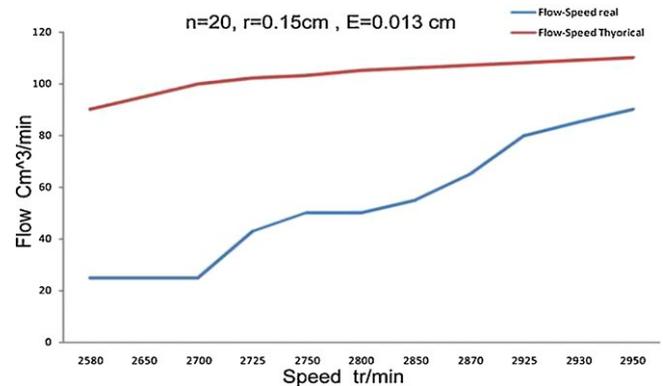


Fig. 9. Preliminary results showing the flow variation versus the rotation speed.

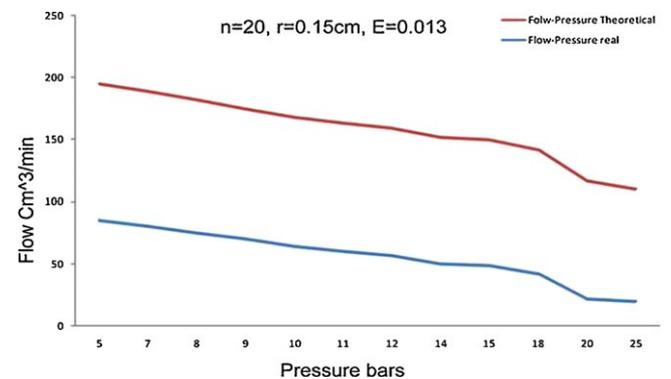


Fig. 10. Preliminary results giving the output pressure versus the produced flow.

## 5. Conclusion

In this paper, the new Integrated Electro-Hydraulic Actuator (IEHA) was modeled base on several subsystems (micro-pump, micro-valve and passive distributor) detailed in the first part of this work. The theoretical input/output relation between the positions of the micro-valve  $X$  and the payload  $Y$  was established using the dynamics of each subsystem. This theoretical modelling leads to a fifth order non-linear differential equation showing our ability to access to several time derivatives of the payload position. A set of simulation for open and closed loop functioning were carried out based on the Matlab-Simulink models developed. This allowed us to establish several interesting results concerning the time response and the needed maximum value of the micro-valve position. In order to validate the concept, several experiments giving the variation of the flow and the pressure function of the eccentricity were carried out using a IEHA prototype and a hydraulic test-bed. Further works concerning the improvements of the IEHA performances are planned, such as pressure, flow and speed are currently under development. The ultimate goal of these developments will be to closely emulate the behaviours of an artificial muscle with intrinsic active compliance properties.

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